

From Perfect Conductor to Perfect Insulator: the Zero-Plateau QH State in Graphene

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Acknowledgements:

L. Brey, R. Berkovits, M. Goldstein, K. Novoselov, P. Ong

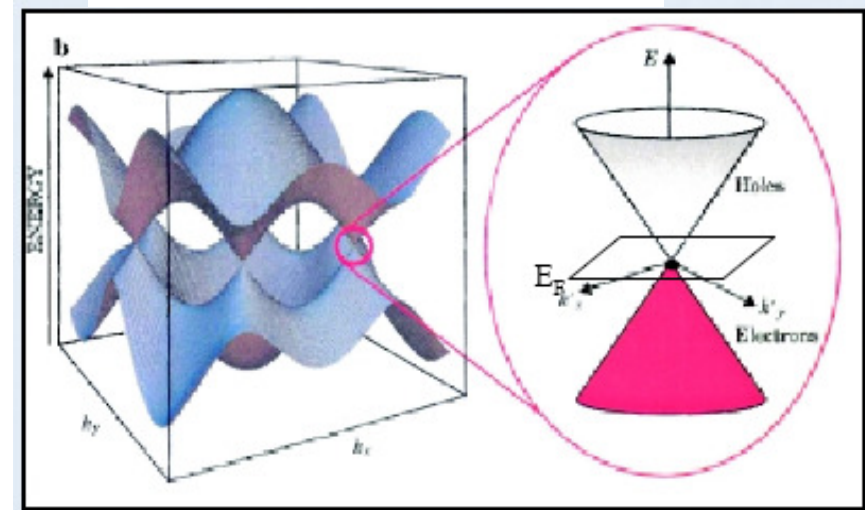
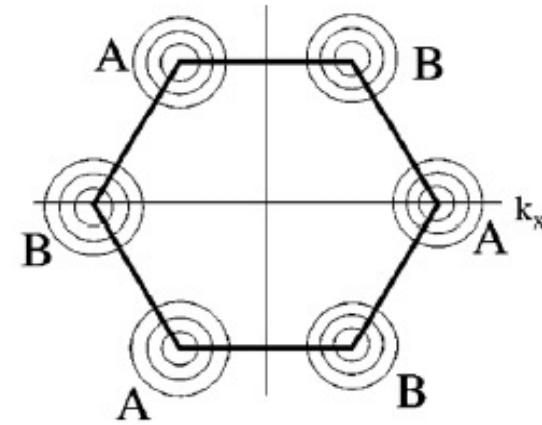
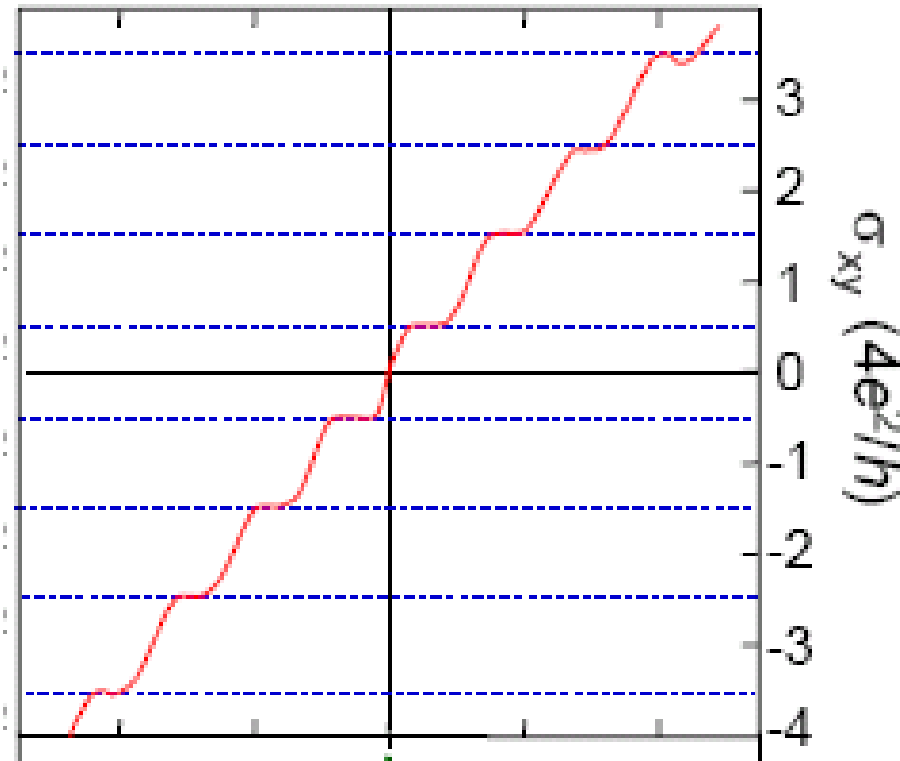
NSF, GIF

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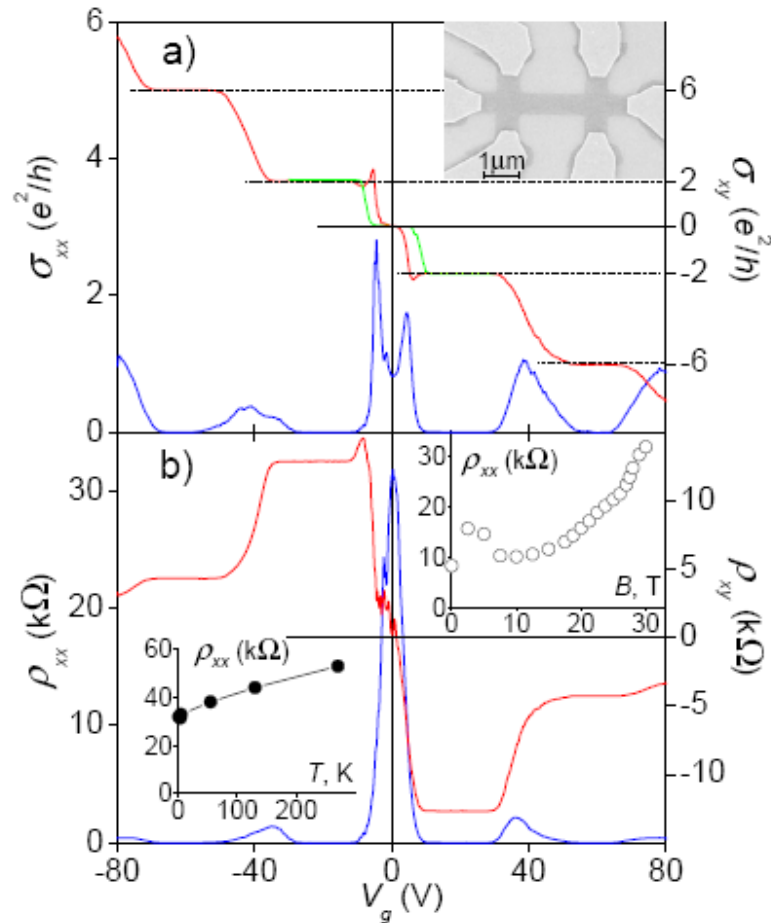
“...Now, **here**, you see, it takes all the running **you** can do, to keep in the same place...” (L. Carroll, from “Alice’s Adventures through the Looking Glass”)

The Quantum Hall Effect in Graphene



Novoselov, Geim et al., Nature **438**, 197 (05)

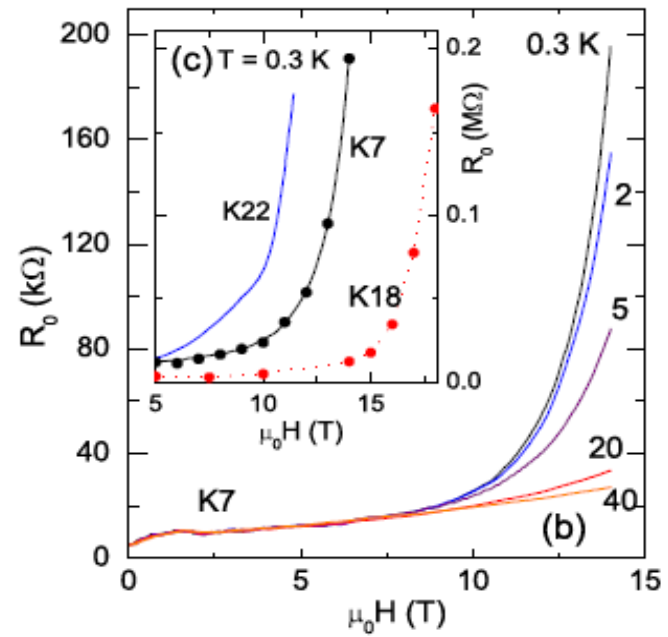
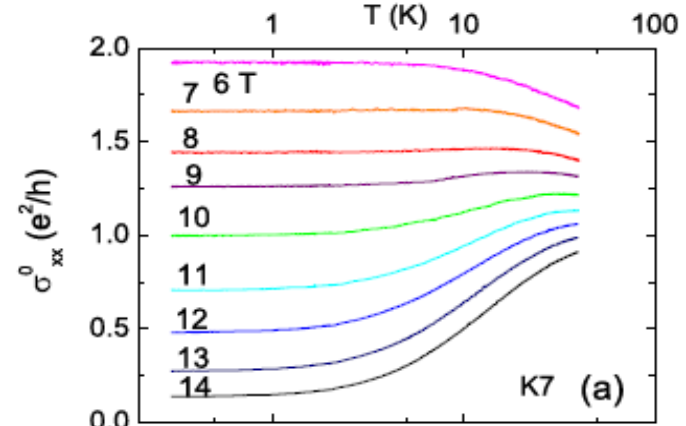
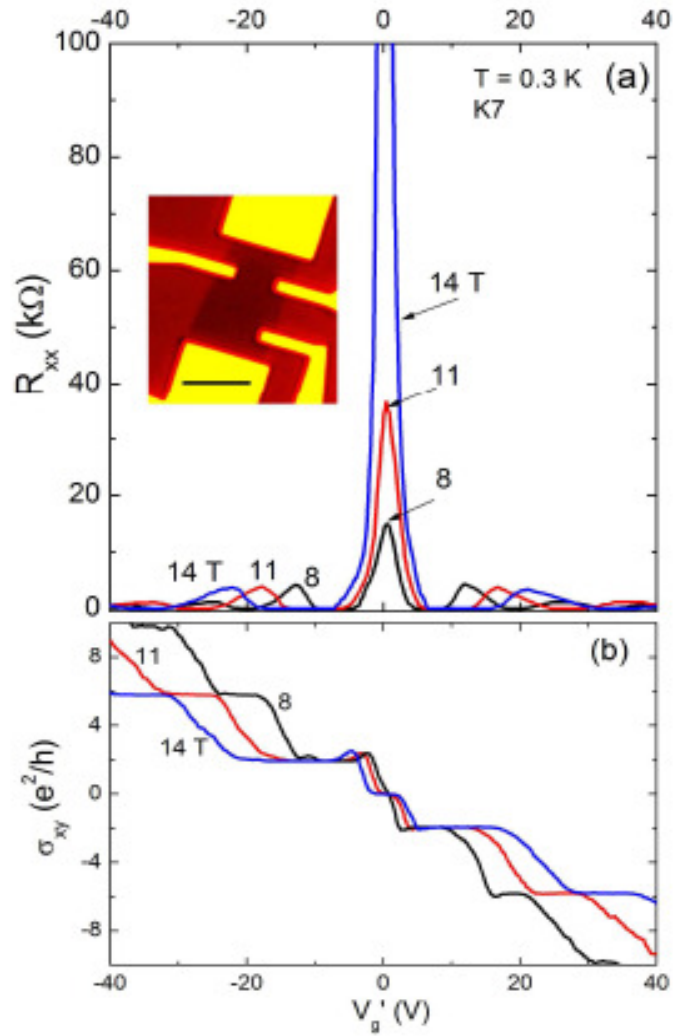
Stronger B: Dissipative QHE at the Dirac Point



Abanin *et al.* (2007)

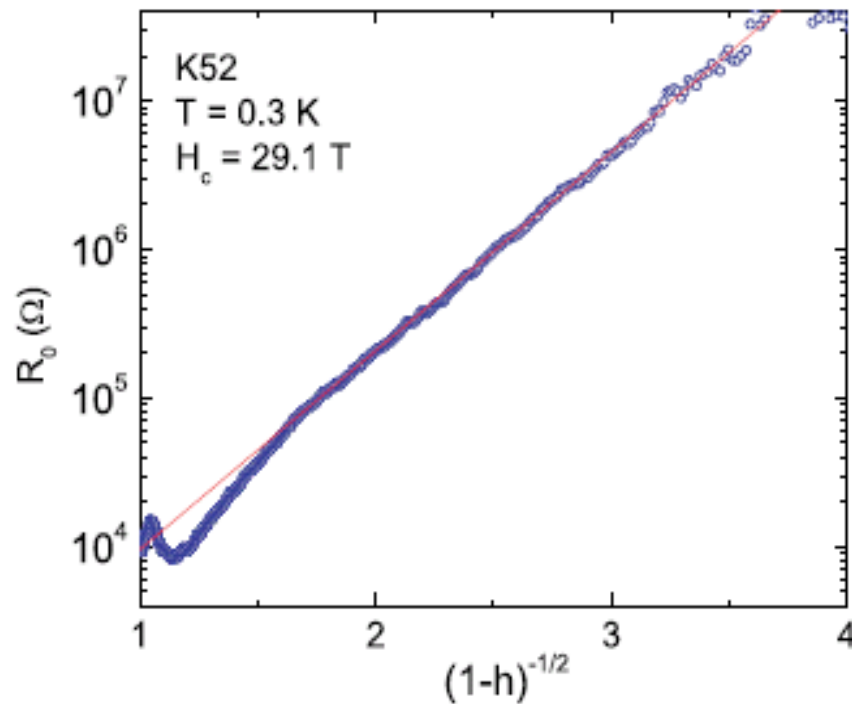
- Plateau at $\nu=0$, but $\rho_{xx} \neq 0$!
- ρ_{xx} increases with $T \Rightarrow$ **Metal**
- ρ_{xx} increases with B

Checkelsky, Li and Ong (2007):



Checkelsky, Li and Ong (2008): more on the “Insulator” -

Divergence of R_0 at high H

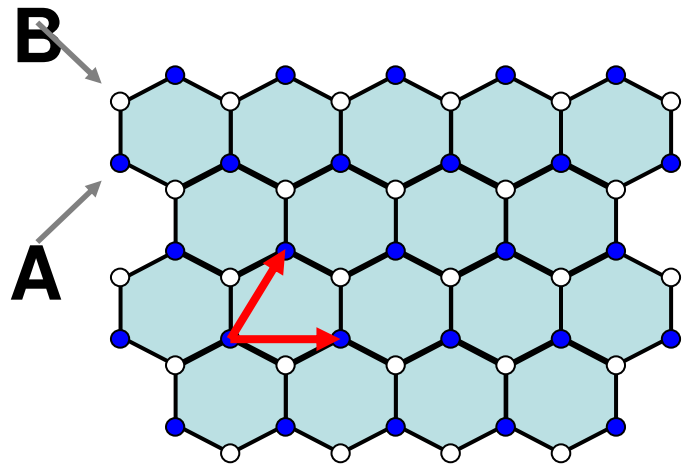


$$R_0 \propto \exp \left\{ \frac{2C}{\sqrt{(H_c - H)}} \right\}$$

$\xi_{KT}^2 !!$

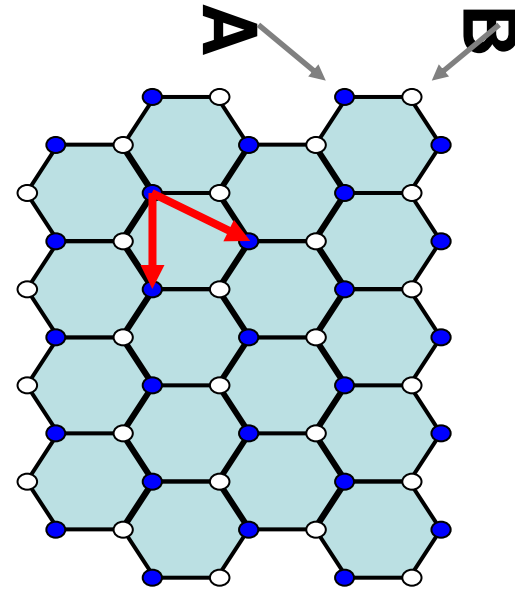
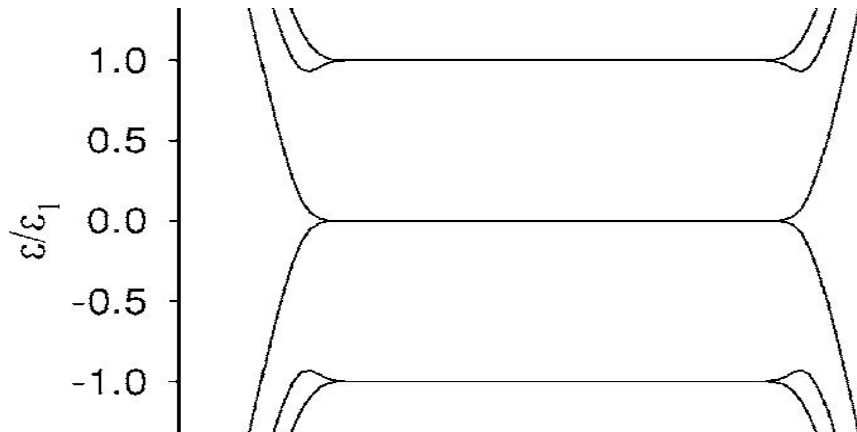
QH Edge-States in Graphene Ribbons

Brey and Fertig (2006):



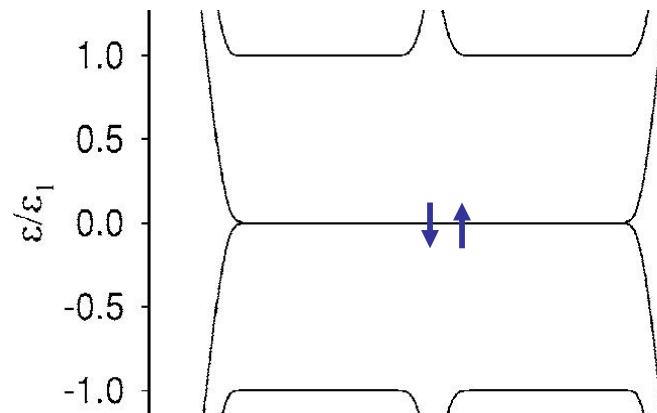
Armchair Edge

Zigzag Edge

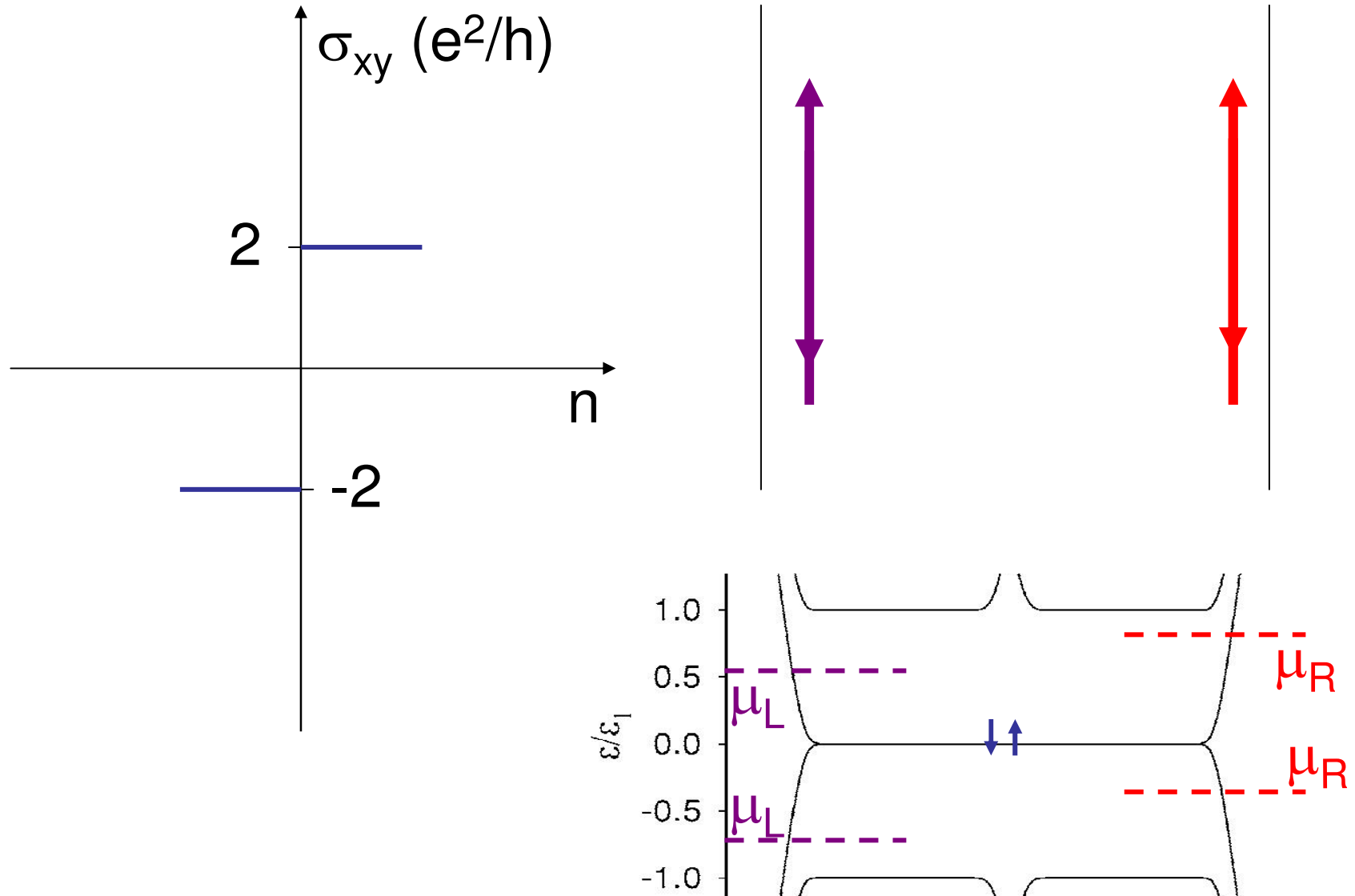


Zigzag Edge

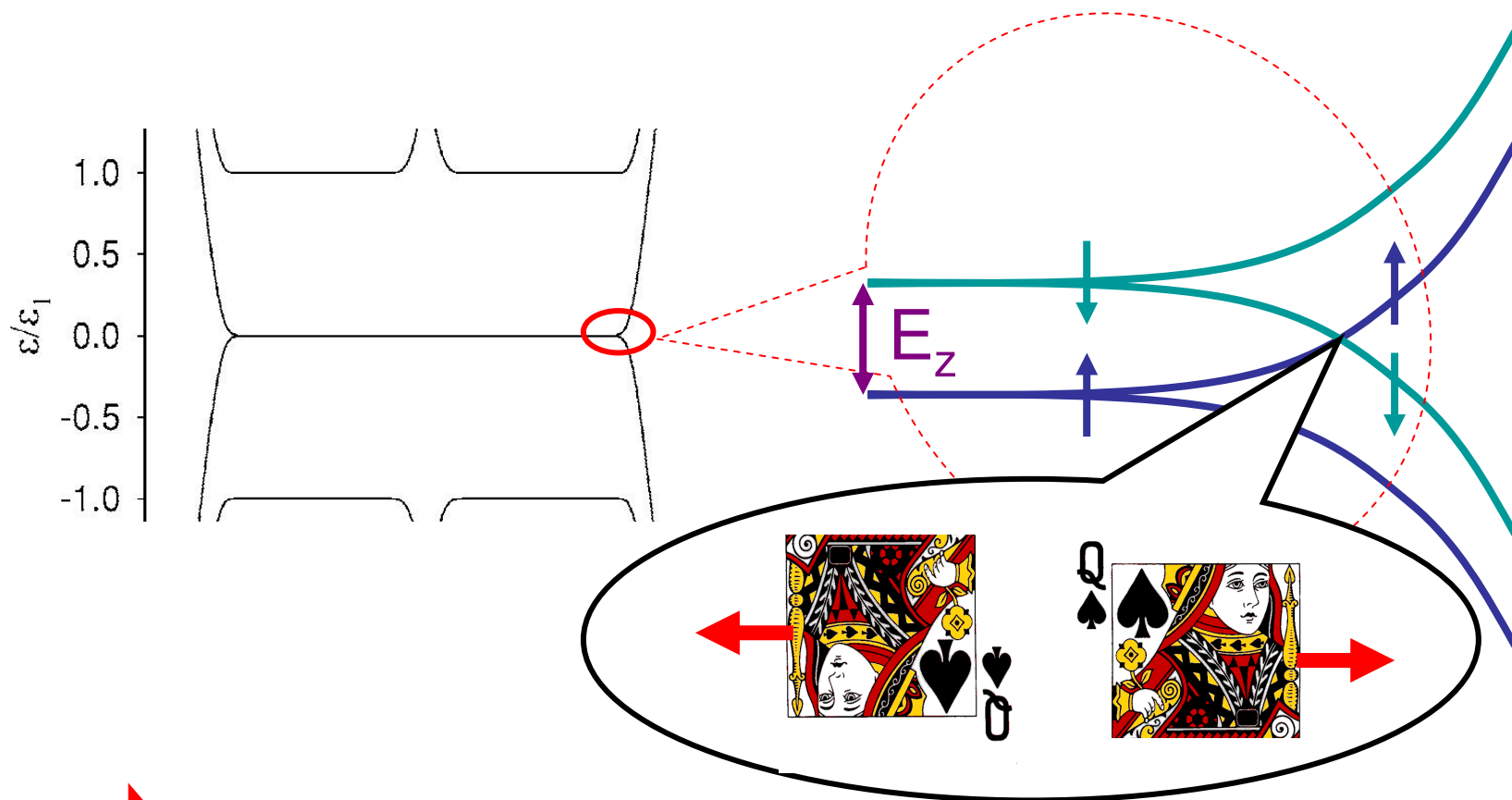
Armchair Edge



QH Edge-States in Graphene Ribbons

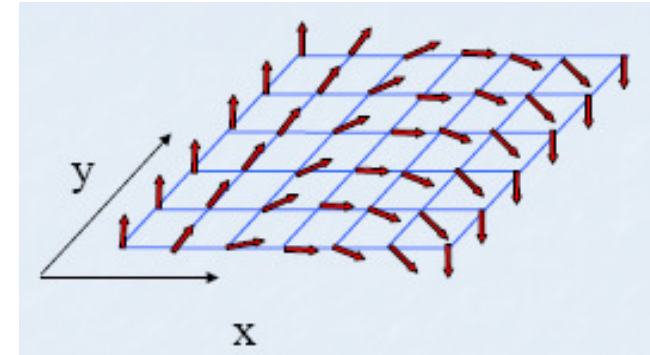
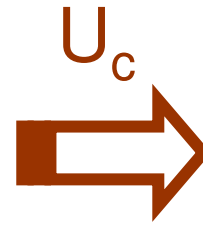
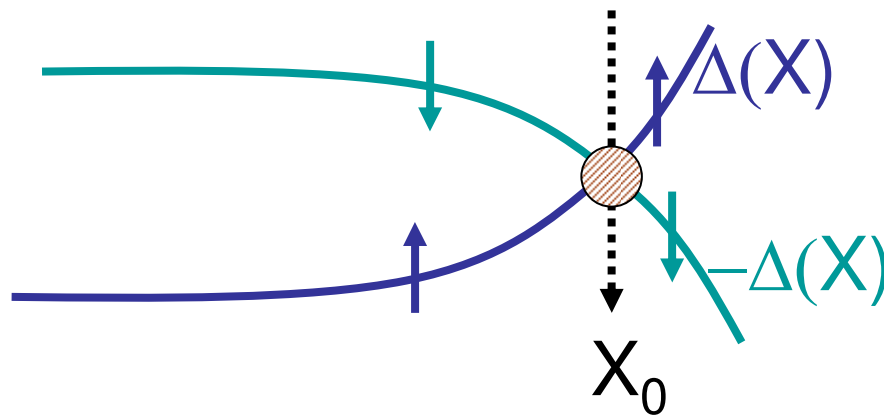


Origin of $\nu=0$ plateau at high H: Zeeman splitting



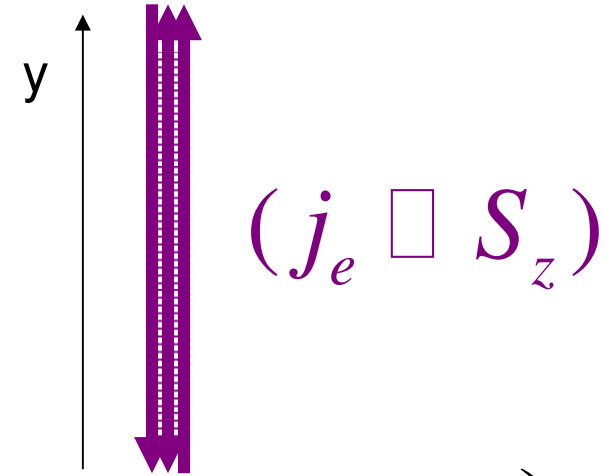
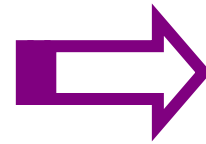
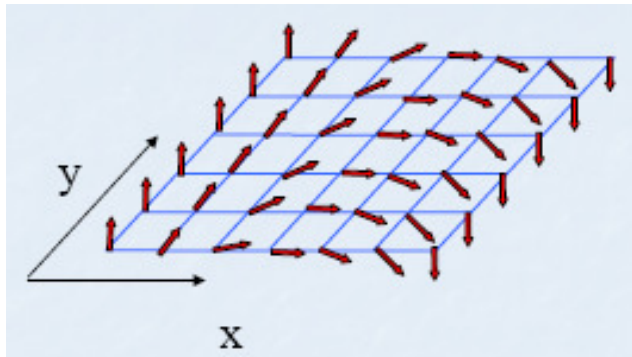
Spin-flip impurities \rightarrow Finite ρ_{xx}
(Abanin, Levitov & Lee)

Role of Coulomb interactions (Fertig & Brey, 2007):



Finite width Domain Wall

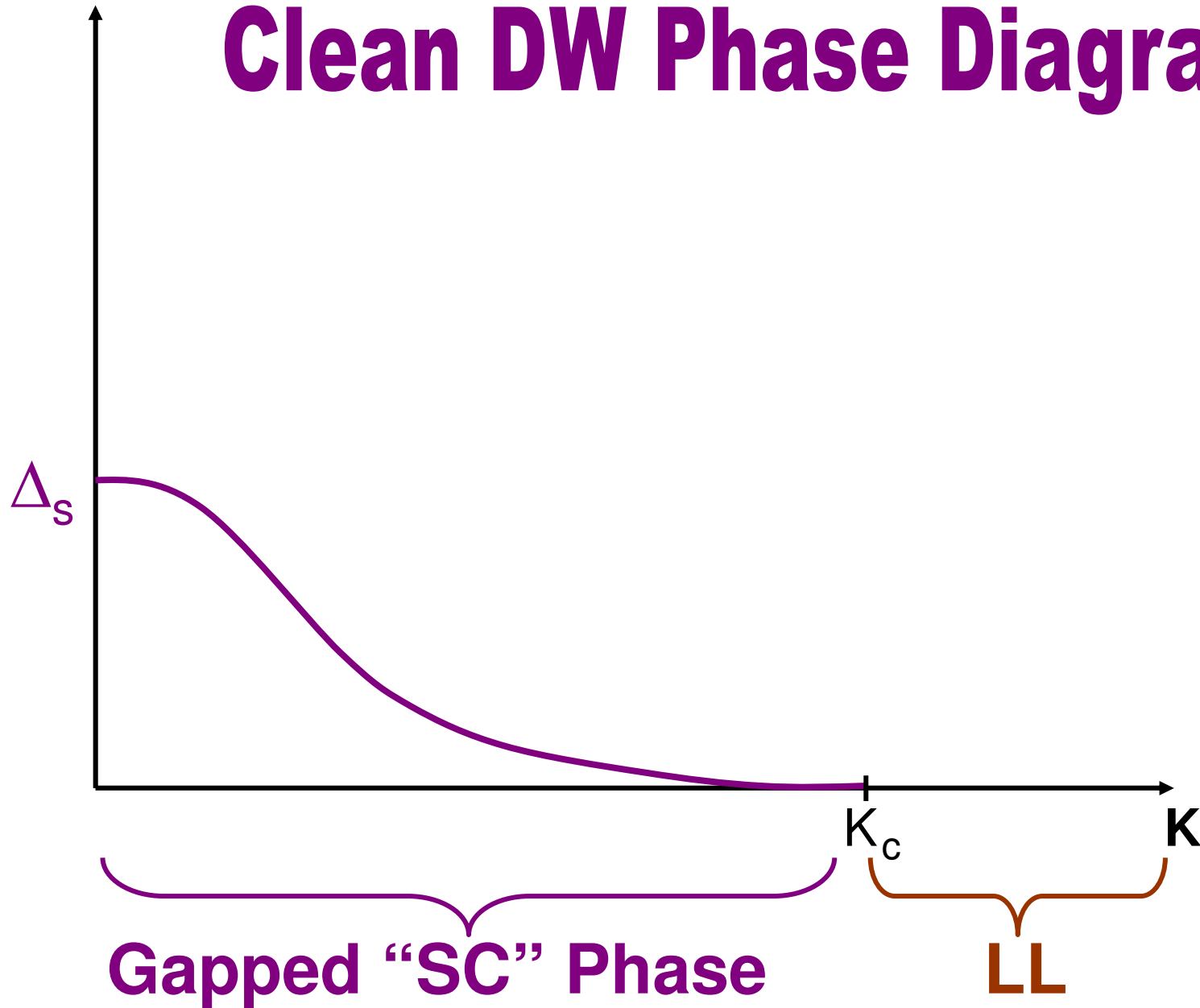
Our Theory: Effective 1D Model for the Domain Wall



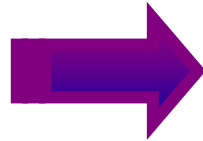
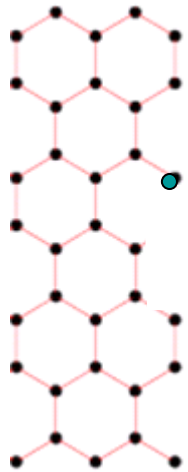
$$H_{DW} = \int \frac{dy}{2\pi} \left(\underbrace{uK (\partial_y \theta)^2}_{\text{"charging energy"}} + \underbrace{\frac{u}{K} (\partial_y \phi)^2}_{\text{"SC stiffness"}} - \boxed{g \cos[4\phi]}_{\text{"Josephson Coupling"}} \right)$$

$K[H, U_c, \Delta'(X_0)]$, $u[H, U_c, \Delta'(X_0)]$, $g[H, U_c, \Delta'(X_0)]$

Clean DW Phase Diagram



Theory for Transport: adding Spin-Flip Interaction



$$H = H_{\text{DW}} + H_{\sigma} + H_{\text{int}}$$

$$H_{\sigma} = \epsilon_z \sigma_z$$

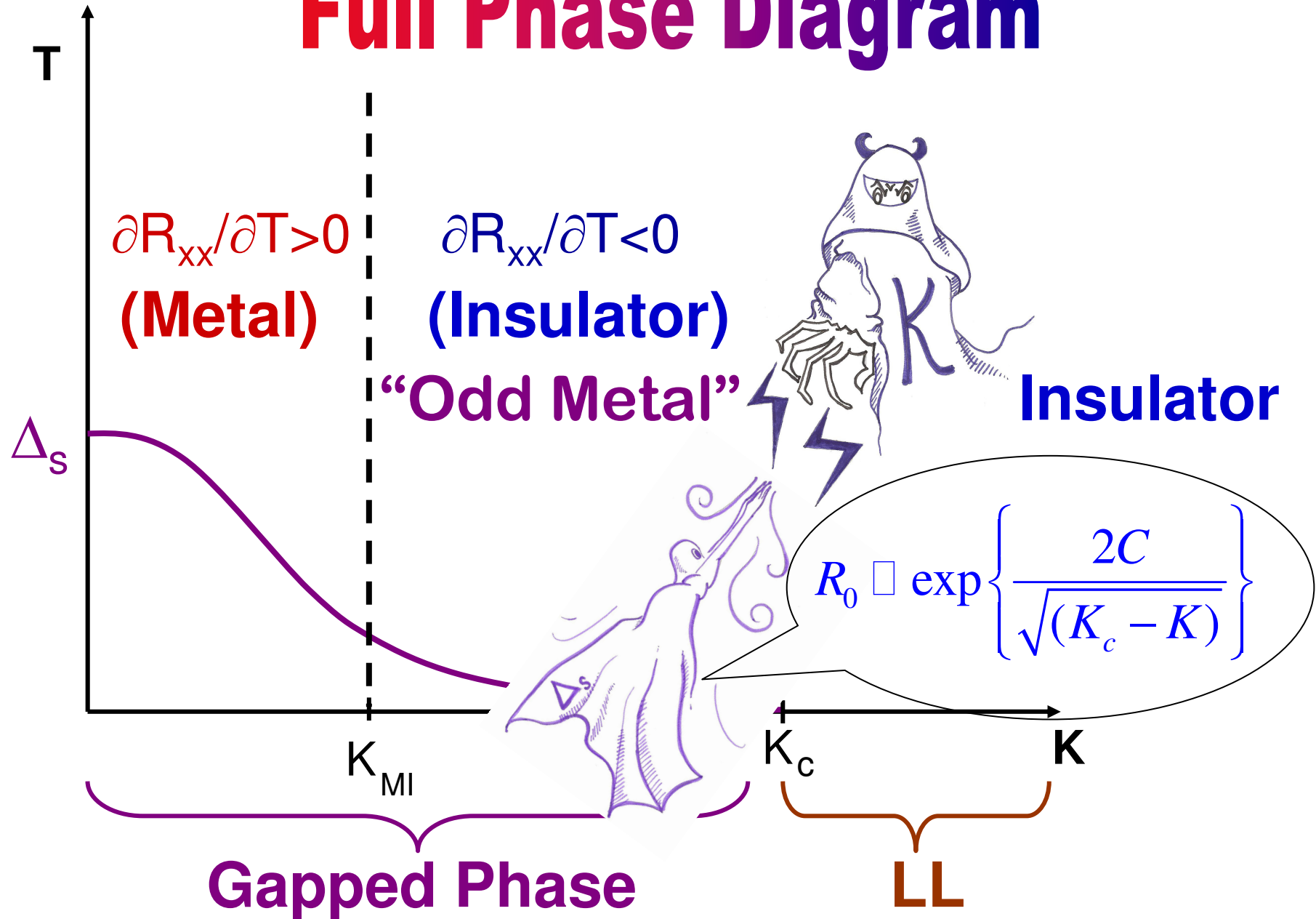
$$H_{\text{int}} = J_k \mathbf{S} \perp \sigma$$

Back-scattering induced resistance for $T > \Delta_s$:

$$R_{xx} \approx g^2 T^{\nu(K)} F[\epsilon_z/T] \quad (g \sim J_k)$$

$$\nu(K_{MI}) = 0$$

Full Phase Diagram



Summary

- ♠ Finite R_{xx} at the $\nu=0$ QH state induced by “chiral Kondo effect”: **Spin-flip** = **Charge backscattering**
- ♠ New type of edge state: **spin Domain Wall** = a non-chiral perfect conducting channel
- ♠ Diverse transport phenomena: $R_{xx}(T)$ is **metallic**, **insulating** or “**odd metal**” depending on K .
Divergence of $R_0 \Rightarrow$ quantum KT-transition (in 1+1d)

